

ROTOR INTERACTION IN THE ANNULUS BILLIARD

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ABSTRACT. Introducing the rotor interaction in the integrable system of the annulus billiard produces a variety of dynamical phenomena, from integrability to ergodicity.

1. INTRODUCTION

The rotor model (see (1.1) below) has been introduced in [MLL] and further studied from the physical point of view in [LLM]. The main motivation for replacing the usual elastic reflection in billiard models with such an interaction is that this way the a priori independent motion of point particles can be coupled. Despite of the physical interest in the subject, there is only a very little amount of mathematically rigorous discussion. This is related to the difficulty of the problem – even the simplest possible *hyperbolic* candidates, planar Sinai billiards with circular scatterers and the related Lorentz gas models turned out to be rather complicated. In [BK] a detailed analysis has been given for a zero measure set of trajectories that possess special symmetries.

In the present article we investigate the rotor interaction in the context of *elliptic* billiards. Consider the motion of point particles in the annulus billiard, i.e. in the planar region bounded by two concentric circles. The one particle system is well known to be integrable ([T] p. 26), while the dynamics of several point particles is the direct product of the individual regular systems. However, if we replace the inner (and possibly the outer) circular billiard scatterer with a rotating object, regularity may be destroyed. Three different models are investigated in this article. In Section 2 we consider the one particle system with the inner circular scatterer rotating and the outer one kept to be a billiard. For this system integrability persists. In Section 3 both scatterers are assumed to rotate and this setting destroys integrability. For a positive measure set of parameters velocity evolution is minimal (Theorem 3.7) and the two dimensional system of our point particle is ergodic (Theorem 3.15). Finally, in Section 4 the motion of *two* point particles is discussed in case the inner scatterer is rotating and the outer one is kept to be a billiard. Though this problem is significantly more complicated, it is possible to find an open set of parameters and, correspondingly, a positive measure set of initial conditions, that this system behaves as

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the one discussed in Section 3 for an arbitrary long finite time (Theorem 4.4 and Remark 4.6). From a physical point of view this result is probably the most interesting as it indicates, as claimed in the physics literature (see [MLL]), the following phenomena: the rotor interaction switches on non-trivial coupling among the particles that would follow independent (in our case integrable) patterns if they were studied in the original billiard setting.

Throughout the paper we consider the motion of point particles in the annular region bounded by two concentric circles, the outer having radius one, the inner radius $R < 1$. Interaction with the circular scatterers is either an elastic reflection (as in billiard models) or the scatterer is assumed to rotate with some angular velocity ω and interact with the point particle according to the laws

$$(1.1) \quad \begin{aligned} v'_n &= -v_n; \\ v'_t &= v_t - \frac{2\eta}{1+\eta}(v_t - r\omega); \\ r\omega' &= r\omega + \frac{2}{1+\eta}(v_t - r\omega). \end{aligned}$$

Here the coordinates v_n and v_t are the normal and tangential components (with respect to the circle of collision) of the particle velocity and the primes indicate post-collisional values. Furthermore r is 1 for the outer and R for the inner circular scatterer, and $\eta = \frac{\Theta}{mr^2}$, where Θ is the moment of inertia for the scatterer in question and m is the mass of the point particle.

Remark 1.1. *Note that for the same velocity vector \vec{v} the splitting into tangential and normal components is different for the inner and the outer scatterer. In Sections 2 and 4 it is only the inner scatterer that is assumed to rotate, and thus the components v_n and v_t in these sections always refer to the splitting at the inner scatterer. In Section 3 both scatterers rotate and it is more convenient to use different notation, see the beginning of Section 3.1.*

Convention 1.2. *The evolution of the coordinates in Formula (1.1) will be referred to as the base.*

In each section the dynamics of the base will be studied before turning to the whole system, which is a skew product over the base, sometimes referred to as fibers.

2. ONE PARTICLE, ONE ROTATING SCATTERER

This section is devoted to a model which is closest to the original integrable annulus billiard, namely the motion of one point particle with only the inner scatterer rotating. For certain initial conditions the point particle never reaches the inner circular scatterer. In this

invariant set motion is that of a point particle in a circle, a system well-known to be integrable. We consider the complement.

At the inner scatterer we have

$$(2.1) \quad \begin{aligned} v'_n &= -v_n; \\ v'_t &= v_t - \frac{2\eta}{1+\eta}(v_t - R\omega); \\ R\omega' &= R\omega + \frac{2}{1+\eta}(v_t - R\omega) \end{aligned}$$

which is just the inner version of Equation (1.1), see Remark 1.1.

Lemma 2.1. (1) *If a trajectory collides once with the inner scatterer, it collides with it infinitely often. More precisely, there is an alternating sequence of collisions, every first with the inner and every second with the outer scatterer.*

(2) *There are three integrals of motion.*

(3) *The velocity motion is periodic of period 2.*

Proof. A trajectory which starts from the inner circle is reflected elastically in the outer circle, thus it returns to the inner circle. Iterating this argument proves statement (1).

Consider two consecutive reflections at the inner scatterer with one billiard reflection at the outer one in between. Let us denote by the triples $v_n(1), v_t(1), \omega(1)$ and $v_n(2), v_t(2), \omega(2)$ the incoming velocities at the two consecutive inner reflections, and let us refer to the outgoing counterparts by primes. By the geometrical picture sketched above:

$$v_n(2) = -v_n(1)'; \quad v_t(2) = v_t(1)'$$

and furthermore $\omega(1)' = \omega(2)$ as the inner scatterer does not interact while the particle collides outside. This, together with (2.1), gives, on the one hand

$$v_n(1) = v_n(2) = -v_n(1)' = -v_n(2)'$$

which can be regarded as an integral of motion. On the other hand, to get $v_t(2)'$ and $\omega(2)'$ we need to iterate the second and the third formulas in (2.1) twice. This gives

$$\omega(2)' = \omega(1); \quad v_t(2)' = v_t(1)$$

thus statement (3) is proved.

Finally multiply the third equation in (2.1) with η and add it to the second to get

$$(2.2) \quad v_t + \eta R\omega = v'_t + \eta R\omega' = N.$$

Furthermore, subtracting the third line from the second in (2.1) we get

$$v'_t - R\omega' = -(v_t - R\omega)$$

which, together with (2.2) results in

$$(2.3) \quad v_t^2 + \eta R^2 \omega^2 = v_t'^2 + \eta R^2 \omega'^2 = E.$$

Thus we have calculated the other two integrals. \square

Remark 2.2. *Equation (2.2) is a multiple of the angular momentum, while (2.3) is the “tangential” part of the full kinetic energy (the normal part of the kinetic energy, i.e. the length of v_n is in itself conserved as the third integral).*

In the space of velocities (v_t, ω) , that can be regarded as \mathbb{R}^2 , fixing the values (2.2) and (2.3) we get a line and a circle, respectively. These two curves intersect in (at most) two points (v_t, ω) and (v_t', ω') . Velocity evolution is a period two cycle on these two points.

Turning to the evolution in the fibers, we essentially get the behavior observed in the integrable annulus billiard. Let us parameterize the outer circle by the arclength parameter φ and denote the points of consecutive impacts on a trajectory by $\varphi_k, k \in \mathbb{Z}$. Then

$$\varphi_{k+1} = \varphi_k + \hat{\beta} + \hat{\beta}'$$

where

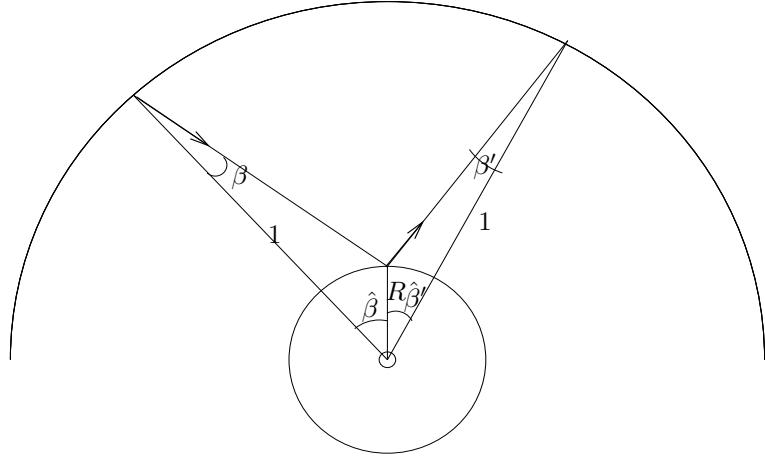
$$\beta = \arctan \frac{v_t}{v_n}, \quad \hat{\beta} = \arcsin \left(\frac{\sin \beta}{R} \right) - \beta$$

and the same for the primed quantities, see Figure 1. Thus the angles $\hat{\beta}, \hat{\beta}' \in [-\arccos(R), \arccos(R)]$ depend smoothly and strictly monotonically on the ratios $\frac{v_t}{v_n}$ and $\frac{v_t'}{v_n'}$, respectively. The consecutive φ_k -s follow an orbit of a circle rotation, which is irrational apart from a countable collection of one codimensional manifolds in the three dimensional parameter space of integrals. To see this keep the values of E and N ((2.3) and (2.2)) fixed and increase the value of v_n continuously, then $\hat{\beta} + \hat{\beta}'$ decreases continuously and is rational only for a countable set of v_n -s.

3. ONE PARTICLE, TWO ROTATING SCATTERERS

3.1. Coordinates, parameters and integrals of motion. In this section we will find that allowing both scatterers to rotate results in a behavior different from the integrable models discussed so far.

As long as velocities are concerned we need to handle an a priori four dimensional system: the coordinates are, on the one hand, ω_1, ω_2 for the angular velocities of the inner and outer scatterers, respectively, and, on the other hand, \vec{v} , the two dimensional velocity vector of the point particle. Recall Remark 1.1: in contrast to the previous section, we will use notations v_n and v_t for the splitting of the velocity \vec{v} into normal and tangential components at the *outer* scatterer. The sign of v_t has a physical meaning, it describes the direction of the “circular”

FIGURE 1. $\hat{\beta} = \arcsin\left(\frac{\sin\beta}{R}\right) - \beta$

component of the motion of the point particle. On the contrary, the sign of v_n is in a certain sense irrelevant, the positivity/negativity of v_n merely tells us if we are just before/after a collision with the outer circle, thus we fix v_n to be always positive. The symbols \bar{v}_n, \bar{v}_t are used for the splitting of \vec{v} into normal and tangential components at the inner scatterer. By simple geometrical observation they can be defined for vectors \vec{v} which satisfy $\frac{v_t^2}{v_n^2} \leq \frac{R^2}{1-R^2}$ as $\bar{v}_t = \frac{1}{R}v_t$ and $\bar{v}_n = \sqrt{v_n^2 - \frac{1-R^2}{R^2}v_t^2}$. Otherwise the trajectory starting from the outer scatterer reaches it again without interacting with the inner one.

We use the notation ω_i for the angular velocity and Θ_i for the moment of inertia with $i = 1, 2$ in the case of the inner and the outer scatterer, respectively. The rescaled moments of inertia are $\eta_1 = \frac{\Theta_1}{mR^2}$ and $\eta_2 = \frac{\Theta_2}{m}$. In these coordinates, the interaction at the outer scatterer is given by:

$$(3.1) \quad \begin{aligned} v'_t &= v_t - \frac{2\eta_2}{1+\eta_2}(v_t - \omega_2); \\ \omega'_2 &= \omega_2 + \frac{2}{1+\eta_2}(v_t - \omega_2). \end{aligned}$$

By the convention on the sign of v_n made above, the normal velocity is preserved and straightforward calculation gives two further preserved quantities:

$$(3.2) \quad v_t + \eta_2\omega_2 = v'_t + \eta_2\omega'_2; \quad v_t^2 + \eta_2\omega_2^2 = v'^2_t + \eta_2\omega'^2_2.$$

As to interaction with the inner scatterer

$$\begin{aligned} \bar{v}'_t &= \bar{v}_t - \frac{2\eta_1}{1+\eta_1}(\bar{v}_t - R\omega_1); \\ R\omega'_1 &= R\omega_1 + \frac{2}{1+\eta_1}(\bar{v}_t - R\omega_1). \end{aligned}$$

Note, however, that $v_t = R\bar{v}_t$ and thus these equations are equivalent to

$$(3.3) \quad \begin{aligned} v'_t &= v_t - \frac{2\eta_1}{1+\eta_1}(v_t - R^2\omega_1); \\ R^2\omega'_1 &= R^2\omega_1 + \frac{2}{1+\eta_1}(v_t - R^2\omega_1). \end{aligned}$$

Two preserved quantities for the inner collisions are

$$(3.4) \quad v_t + \eta_1 R^2 \omega_1 = v'_t + \eta_1 R^2 \omega'_1; \quad v_t^2 + \eta_1 R^4 \omega_1^2 = v'^2_t + \eta_1 R^4 \omega'_1^2.$$

As ω_1 (ω_2) does not change at outer (inner) collisions, (3.2) and (3.4) imply the presence of two integrals of motion, the angular momentum

$$(3.5) \quad v_t + \eta_1 R^2 \omega_1 + \eta_2 \omega_2 = \text{const.} = N$$

and the tangential energy

$$(3.6) \quad v_t^2 + \eta_1 R^4 \omega_1^2 + \eta_2 \omega_2^2 = \text{const.} = E(> 0).$$

The term “tangential” refers to the fact that the full kinetic energy of the system is a third integral of motion, always greater than the tangential one. More precisely, dividing the full energy $\frac{1}{2}(m\vec{v}^2 + \Theta_1\omega_1^2 + \Theta_2\omega_2^2)$ by $\frac{1}{2}m$ we get

$$(3.7) \quad v_n^2 + v_t^2 + \eta_1 R^2 \omega_1^2 + \eta_2 \omega_2^2 = \text{const.} = F(> E).$$

Note that (3.2) ((3.4)) and (3.7) together imply that v_n (\bar{v}_n) is preserved during inner (outer) collisions.

Remark 3.1. (1) One of our aims below is to understand how the dynamical behavior depends on the several parameters of the system. The constants η_1 and η_2 will be referred to as physical parameters in contrast to the integrals of motion N, E and F .

(2) To avoid complications we assume $\vec{v} \neq \vec{0}$. Note that this assumption is reasonable in the following sense: $\vec{v} = \vec{0}$ is only possible for a codimension one submanifold of integrals of motion. To see this observe that if we had $\vec{v} = \vec{0}$, (3.5) and (3.6) would fix the values of the other two coordinates ω_1 and ω_2 (up to two possibilities) and thus would create, via (3.7), a relation between the three integrals of motion.

3.2. Evolution of velocities. For convenience we rescale our velocities as

$$(3.8) \quad x := \sqrt{\eta_1} R^2 \omega_1, \quad y := \sqrt{\eta_2} \omega_2, \quad z := v_t, \quad w := v_n.$$

In these rescaled quantities our integrals of motion are

$$(3.9) \quad x^2 + y^2 + z^2 = E; \quad \sqrt{\eta_1}x + \sqrt{\eta_2}y + z = N$$

and

$$(3.10) \quad \frac{x^2}{R^2} + y^2 + z^2 + w^2 = F.$$

By (3.1) and (3.3) we can first consider how the velocity coordinates x, y, z evolve at collisions. With the convention on its sign the value of $w (= v_n)$ is determined by (3.10) for each fixed value of the other three coordinates. In the rest of the section, unless otherwise stated, velocity evolution will always refer to the dynamics in the triple of x, y, z .

Taking into account that x (y) does not change at outer (inner) collisions, in the rescaled coordinates the transformations (3.1) and (3.3) are reflections across the planes

$$(3.11) \quad \sqrt{\eta_2}z - y = 0 \quad \text{and} \quad \sqrt{\eta_1}z - x = 0,$$

respectively. Furthermore, by (3.9) motion is restricted to the intersection of a sphere and a plane, i.e. to a circle which will be referred to as the velocity circle. We will see that the base, see Convention 1.2, will consist of two copies of this circle. Note that these circles do not describe the physical position of the particle, on the other hand they give the state of the velocities. On the velocity circle the reflections across the planes (3.11) reduce to reflections across lines. For the outer collisions we have reflections across

$$(3.12) \quad l_2(t) = \left(-\frac{1 + \eta_2}{\sqrt{\eta_1}}t, \sqrt{\eta_2}t, N + t \right)$$

while for the inner ones across

$$l_1(t) = \left(\sqrt{\eta_1}t, -\frac{1 + \eta_1}{\sqrt{\eta_2}}t, N + t \right).$$

Let us denote the angle between these two lines by $\frac{\gamma}{2}$. Depending on it we distinguish rational, irrational (and later on even Diophantine) cases.

Lemma 3.2. *Rationality is independent of the integrals of motion. On the other hand, there exists a countable set of codimension 1 submanifolds of physical parameters which lead to the rational case, otherwise (and thus for a full Lebesgue measure set of parameters) γ is irrational.*

Proof. Straightforward calculation gives

$$(3.13) \quad \cos \frac{\gamma}{2} = ((1 + \eta_1^{-1})(1 + \eta_2^{-1}))^{-\frac{1}{2}}.$$

Thus the rationality of γ depends only on the physical parameters. Consider the numbers $\cos(\frac{p\pi}{q})$ for positive integers $p < q$. Note that elements of this countable set A are all algebraic numbers. The base is rational iff the left hand side of (3.13) is equal to some number in A . This happens for a countable collection of one codimensional submanifolds of physical parameters. \square

Convention 3.3. *To agree with different standard conventions, in the geometric framework we will think of \mathbb{S}^1 as $[0, 2\pi]$, while in the number*

theoretic framework we think of \mathbb{S}^1 as $[0, 1]$. In both cases we scale all \mathbb{S}^1 valued quantities $(s, \varphi, \gamma, \alpha(s))$ appropriately. This should cause no confusion.

To give a description of the base dynamics we introduce the arclength parameter s as a coordinate on the velocity circle. For definiteness $s = 0$ corresponds to one of the intersection points of the circle with the line (3.12). In this coordinate the outer and inner bounces, i.e. the transformations (3.1) and (3.3) are

$$(3.14) \quad s \rightarrow s_1 = -s \pmod{1} \text{ and}$$

$$(3.15) \quad s \rightarrow s_2 = -s - \gamma \pmod{1},$$

respectively. Take two copies, O and I , of the velocity circle describing the outgoing velocities just after the bounces at the inner and outer circles, respectively. The phase space of the velocity motion, or the base, is

$$\hat{M} = \mathbb{S}^1 \times \{1, 2\} = I \cup O; \quad \hat{M} \ni \bar{s} = (s, i) \text{ with } s \in \mathbb{S}^1 \text{ and } i = 1, 2.$$

The velocity evolution will be denoted by $T : \hat{M} \rightarrow \hat{M}$. As an inner bounce is always followed by an outer one, for $\bar{s} \in I$, i.e. for $(s, 1)$ the image $T\bar{s}$ is always in O , more precisely $T(s, 1) = (s_1, 2)$ (recall (3.14)). An outer bounce, however, may be followed by either an inner or an outer bounce, depending on the angle the point particle velocity makes with the normal vector of the outer scatterer.

Lemma 3.4. (1) There is a (possibly empty) open set $U \subset \mathbb{S}^1$, such that for $\bar{s} \in U$ (more precisely for $\bar{s} = (s, 2)$ with $s \in U$), we have $T\bar{s} = (s_1, 2)$ (the image is the point defined by (3.14) in O), and for $\bar{s} \in O \setminus U$ (more precisely for $\bar{s} = (s, 2)$ with $s \in \mathbb{S}^1 \setminus U$) we have $T\bar{s} = (s_2, 1)$ (the image is the point defined by (3.15) in I).

(2) The set $U \subset \mathbb{S}^1$ is invariant with respect to the reflection (3.15).

(3) There is an open set of integrals of motion for which U is empty.

Remark 3.5. We will refer to empty U as the alternating case, non-empty U as the non-alternating case. Both of them occur for open sets of integrals of motion, independently of the value of physical parameters.

Proof of Lemma 3.4.

(1) By Formulas (3.9) and (3.10) the arclength s on the velocity circle determines $z = v_t$ and $w = v_n \geq 0$ continuously. The geometric condition for $(s, 2) = \bar{s} \in O$ to reach the outer scatterer without touching the inner one (and thus interacting with it) is $\frac{v_t^2}{v_n^2} > \frac{R^2}{1-R^2}$. Note that at tangential collisions with the inner scatterer v_t is modified by (3.3), thus this corresponds to landing on I .

We define

$$(3.16) \quad U := \left\{ s \in \mathbb{S}^1 \mid z^2(s) > \frac{R^2}{1-R^2} w^2(s) \right\}.$$

If $\bar{s} \in O$ with $s \in U$ (i.e. in case the next bounce is outside) then (3.14) applies, and thus $T\bar{s} = (s_1, 2)$. Otherwise (i.e. if $\bar{s} \in O \setminus (U \times \{2\})$) the next bounce is inside, thus the image is the point defined by (3.15) in I , or $T\bar{s} = (s_2, 1)$. The set $U \subset \mathbb{S}^1$ is the preimage of an open set by a continuous map, thus it is open.

(2) To see that U is invariant under (3.15) note first that y does not change during an inner collision. This, together with the presence of the integrals (3.9) and (3.10) implies that $w^2 - \frac{1-R^2}{R^2} z^2$ is invariant for inner bounces, i.e. for the reflections (3.15). The statement follows as it is exactly the sign of this quantity that distinguishes the points of U .

(3) The set U is empty if

$$(3.17) \quad z^2(s) < \frac{R^2}{1-R^2} w^2(s)$$

for all s . Consider the integrals of motion defined by (3.9) and (3.10) and denote $\lambda = \frac{F}{E} > 1$. By (3.9) x^2 and z^2 are both less than or equal to E for all s . This, together with (3.10) gives

$$\begin{aligned} w^2 &\geq F - y^2 - z^2 - x^2 - x^2(R^{-2} - 1) \geq F - E - (R^{-2} - 1)E = \\ &= F - ER^{-2} = E(\lambda - R^{-2}) \geq (\lambda - R^{-2})z^2 \end{aligned}$$

for all s . This implies U is empty if $\lambda - R^{-2} > \frac{1-R^2}{R^2}$, i.e. if $\lambda > (\frac{2}{R^2} - 1)$, a condition satisfied for a set of parameters that has nonempty interior. \square

Remark 3.6. (1) Actually U is empty iff we replace the strict inequality in Formula (3.17) with a nonstrict one. The analysis of the equality case is substantially more complicated, and is a co-dimension one phenomenon in the set of parameters, thus we do not analyze it.

(2) Note that the points $s \in U$ correspond to velocity configurations that cannot be realized at the inner circle. Thus for nonempty U the phase space is $M \subset \hat{M}$:

$$M = I \cup O = ((\mathbb{S}^1 \setminus U) \times \{1\}) \cup (\mathbb{S}^1 \times \{2\})$$

and the dynamics acts as

$$\begin{aligned} T(O \setminus (U \times \{2\})) &= (\mathbb{S}^1 \setminus U) \times \{1\} = I; \\ T(U \times \{2\}) \cup TI &= \mathbb{S}^1 \times \{2\} = O; \end{aligned}$$

with

$$(3.18) \quad \begin{aligned} T(s, 2) &= (s_2, 1) \quad \text{for } s \in \mathbb{S}^1 \setminus U; \\ T(s, 2) &= (s_1, 2) \quad \text{for } s \in U; \text{ and} \\ T(s, 1) &= (s_1, 2) \quad \text{for } s \in \mathbb{S}^1 \setminus U; \end{aligned}$$

which is well-defined and invertible by statement (2) of Lemma 3.4. As U has no inside copy sometimes we refer to $U \times \{2\}$ as U for brevity, this should cause no confusion.

The main result on the base dynamics is the following theorem.

Theorem 3.7. *In the alternating irrational case the dynamics of the base is minimal, while in all the other cases velocity motion is always periodic. Moreover, all possible periods are even with length bounded above by a constant (which depends on the parameters).*

Proof. Consider the alternating case first. For $\bar{s} \in O$ all odd iterates are in I and all even iterates are in O (and for $\bar{s} \in I$ the other way round). Thus it is enough to consider the first return map onto O (which is exactly T^2). By (3.14) and (3.15) this is a rotation of \mathbb{S}^1 by γ , which is periodic/minimal exactly in the rational/irrational cases. (The odd iterates make the reflected image of the same rotation orbit in I .)

Turning to the non-alternating case, i.e. to $U \neq \emptyset$, first note that $U \cap T^{-1}U$ is an invariant set. To see this consider $(s, 2) = \bar{s} \in U \cap T^{-1}U$. Then both $T\bar{s}$ and $T^2\bar{s}$ are in O . More precisely

$$T(s, 2) = (s_1, 2) \quad \text{and} \quad T^2(s, 2) = T(s_1, 2) = (s, 2),$$

thus the set is invariant, furthermore, all points of it are periodic of period 2.

Now take $(s, 2) = \bar{s} \in U \setminus T^{-1}U$. Then $T\bar{s}$ is in $O \setminus U$, $T^2\bar{s} \in I$ and even $T^{-1}\bar{s} \in I$, as the point is assumed not to belong to the above described invariant set. Moreover, by (3.14) and (3.18)

$$(3.19) \quad T\bar{s} = T(s, 2) = (s_1, 2) \quad \text{and} \quad T^{-1}\bar{s} = T^{-1}(s, 2) = (s_1, 1),$$

thus the trajectory essentially ‘turns back’. The time reflection symmetry of (3.19) is preserved until the next U -hitting time. To see this let $k > 1$ be the smallest integer (possibly ∞) such that $T^k\bar{s} \in U$. Up to time k the orbit is alternating, i.e. by induction

$$\begin{aligned} T^{2l-1}\bar{s} &= (s', 2), & T^{2l}\bar{s} &= (s'_2, 1), \\ T^{-(2l-1)}\bar{s} &= (s', 1), & T^{-2l}\bar{s} &= (s'_2, 2) \end{aligned}$$

for some $s' \in \mathbb{S}^1 \setminus U$ whenever $0 < 2l < k$. In particular $T^{(k-1)}\bar{s} = (s'', 1)$ and $T^{-(k-1)}\bar{s} = (s'', 2)$ for some $s'' \in T^{-1}U \setminus U$ (note that k is always odd). This implies by (3.18) that the preimage of $T^{-(k-1)}\bar{s}$ is not in I , alternation ceases and

$$T^{-k}\bar{s} = (s''_1, 2) \in U \subset O.$$

Moreover

$$T^k \bar{s} = (s''_1, 2) = T^{-k} \bar{s},$$

thus the point is periodic with period $2k$. This way we have shown that any orbit which enters U twice is periodic.

Consider any $\bar{s} = (s, 2) \in O$ and the subset of \mathbb{S}^1 generated from s by the reflections (3.14) and (3.15). This subset of \mathbb{S}^1 consists of two rotation orbits: the first by γ and the second by $-\gamma$, which is the reflected image of the first one by (3.15).

In the non-alternating rational case there are two possibilities. On the one hand, if both of the above mentioned two rotation orbits avoid U , dynamics is that of the alternating case: \bar{s} is periodic of length that depends only on γ (i.e. on the physical parameters). On the other hand, by the invariance of U with respect to (3.15) (cf. statement (2) of Lemma 3.4) if one of the rotation orbits hit U , the other orbit – the reflected image – should hit it as well. Thus if this occurs the trajectory of \bar{s} enters $U \times \{2\}$ twice and becomes periodic with even length. Since the longest possible period occurs when we avoid U , the length of all possible periods is bounded above by a constant that depends only on γ .

In the non-alternating irrational case, as both rotation orbits are dense and thus cannot avoid U , the second of the two possibilities sketched above occurs for all \bar{s} . To complete the proof of the theorem we only need to show that the length of possible periods is bounded above. By minimality and compactness there exists a positive integer N such that the first N rotation-preimages of U cover \mathbb{S}^1 . This implies that the length of any period is less than or equal to $4N$. Moreover, N depends only on γ and U , i.e. on the physical parameters and the integrals of motion. \square

3.3. Analysis of the skew product. In this subsection we extend our investigation to the fibers, i.e. to the question how the points of impact for the consecutive outer bounces follow each other on the outer scatterer. Let us parameterize the outer circle by the arclength parameter φ (the starting point is irrelevant and may be chosen arbitrarily). We consider the skew product

$$(3.20) \quad F : \mathbb{S}^1 \times \mathbb{S}^1 \rightarrow \mathbb{S}^1 \times \mathbb{S}^1; \quad F(s, \varphi) = (T_O s, \varphi + \alpha(s)).$$

Here T_O is the (projection onto \mathbb{S}^1 of the) first return map onto O for T from Section 3.2. Note that T_O is either the first or the second iterate of T . More precisely, if $s \in U$, we have $T_O = T$ and $\alpha(s) = \pi - 2\beta(s)$ where $\tan \beta(s) = \frac{v_t(s)}{v_n(s)} = \frac{z(s)}{w(s)}$. On the other hand, if $s \in \mathbb{S}^1 \setminus U$, we have $T_O = T^2$ and $\alpha(s)$ can be calculated using the geometry of Figure 1

(cf. the end of Section 2) as:

$$(3.21) \quad \begin{aligned} \alpha(s) &= \hat{\beta}(s) + \hat{\beta}(s_2) \quad \text{where} \\ \sin(\beta(s) + \hat{\beta}(s)) &= \frac{\sin \beta(s)}{R}, \\ \tan \beta(s) &= \frac{v_t(s)}{v_n(s)} = \frac{z(s)}{w(s)} \end{aligned}$$

and s_2 is defined by (3.15). Note that, when restricted to U or $O \setminus U$, both α and T_O are continuous.

Recall from Theorem 3.7 that apart from the alternating irrational case T_O is periodic for any s with some period $N = N(s)$. In these periodic situations

$$F^N(s, \varphi) = (s, \varphi + \alpha_N(s)); \quad \alpha_N(s) = \alpha(s) + \alpha(T_O s) + \cdots + \alpha(T_O^{N-1} s).$$

Thus a suitable iterate of the dynamics in the fibers is a rotation, the angle of which depends only on the initial velocity s , and not on the initial position φ . The iterates $F^{k_0+kN}s$, with $0 < k_0 < N$ fixed and $k \in \mathbb{Z}$, are the shifted versions of the same rotation orbit. Moreover, by piecewise continuity of α and T_O , and by the uniform upper bound on $N(s)$, the rotation angle $\alpha_N(s)$ depends on s in a piecewise continuous manner as well (at discontinuity points, however, even $N(s)$ may jump).

Remark 3.8. *Note that in the piecewise continuity statements we have finitely many pieces since the set U , being defined by algebraic equations and inequalities, has finitely many connected components.*

From this point on we restrict to the alternating irrational case. Thus T_O is an irrational rotation, and we have

$$(3.22) \quad F(s, \varphi) = (s + \gamma, \varphi + \alpha(s))$$

with γ irrational and $\alpha : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ continuous. Such systems are called Anzai skew products after their first investigator ([A]). There is an extensive literature on the subject, see [F], [BM], [I], [ILM] and references therein. Some simple issues are even treated in general monographs on dynamical systems like [KH].

In addition to the skew product map (3.22) we are interested in the dynamical behavior in real time, i.e. in a suspension flow over our skew product. The ceiling function for this flow, i.e. the physical (real-valued) first return time to the outer scatterer depends only on the initial velocity, thus it may be denoted as $\tau(s)$ with $\tau : \mathbb{S}^1 \rightarrow \mathbb{R}^+$. We

may calculate $\tau(s)$ by Figure 2 as

$$(3.23) \quad \begin{aligned} \tau(s) &= \frac{d(s)}{|\vec{v}(s)|} + \frac{d(s_2)}{|\vec{v}(s_2)|} \quad \text{with} \\ |\vec{v}(s)| &= \sqrt{v_t(s)^2 + v_n(s)^2} = \sqrt{z(s)^2 + w(s)^2} \quad \text{and} \\ d(\beta(s)) &= d(s) = \cos \beta(s) - \sqrt{R^2 - \sin^2 \beta(s)}, \quad \text{where} \\ \tan \beta(s) &= \frac{v_t(s)}{v_n(s)} = \frac{z(s)}{w(s)} \end{aligned}$$

and s_2 is defined by (3.15).

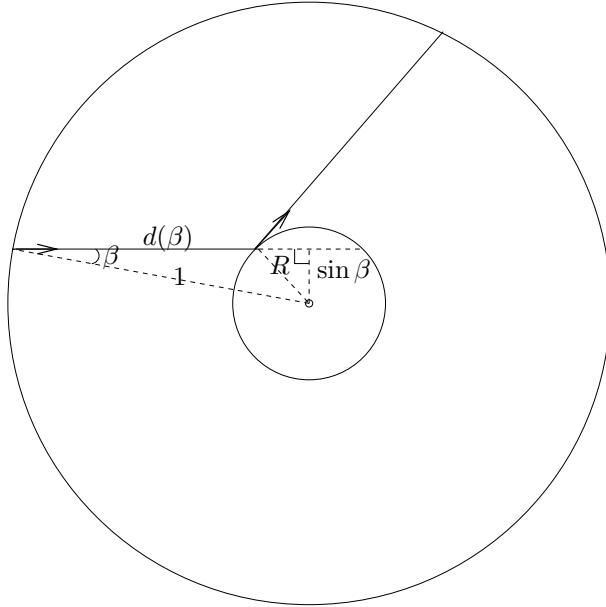


FIGURE 2. $d(\beta) = \cos \beta - \sqrt{R^2 - \sin^2 \beta}$

We recall some basic concepts and properties related to irrational numbers γ and continuous functions $\alpha : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ (on details see [KH] and [B]).

Definition 3.9. *Given $r > 1$ an irrational number $\gamma \in [0, 1]$ is **Diophantine of type r** if there exists $c > 0$ such that for any $p, q \in \mathbb{Z}$, $q \neq 0$*

$$|q\gamma - p| > cq^{-r}.$$

A number which satisfies the above condition for some $r > 1$ is called Diophantine. Liouville numbers γ are those irrationals that do not satisfy the Diophantine condition for any r . It is not hard to see that the set of Diophantine numbers of type r is, as a subset of $[0, 1]$, on the one hand, of the first category, and, on the other hand, of full Lebesgue measure, with, even more, its complement having Hausdorff dimension r^{-1} . Consequently, the set of Diophantine numbers is of full measure

and first category, while its complement, the union of Liouville numbers and rational ones, has zero Hausdorff dimension.

Definition 3.10. For any continuous map $\alpha : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ there exist countably many continuous functions $A : \mathbb{R} \rightarrow \mathbb{R}$ such that $A(s+1) - A(s)$ is an integer independent of s and $\alpha(s)$ can be regarded as the restriction of $A(s)$ to $[0, 1]$. The integer $A(s+1) - A(s)$ is independent of the lift and is called **the degree** of $\alpha(s)$. The map $\alpha(s)$ can be regarded as a real-valued continuous function defined on \mathbb{S}^1 iff its degree is zero.

Definition 3.11. Two measurable maps $\alpha : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ and $\alpha' : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ are said to be **cohomologous** (for the rotation by γ) iff there exists $V : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ measurable such that

$$\alpha(s) = \alpha'(s) + V(s + \gamma) - V(s)$$

for almost all $s \in \mathbb{S}^1$. Maps cohomologous to 0 are called coboundaries.

The following statements date back to [A].

Lemma 3.12. (1) The skew product (3.22) is ergodic (with respect to Lebesgue measure) iff there is no $p \in \mathbb{Z} \setminus \{0\}$ such that $p\alpha(s)$ is a coboundary.

(2) The skew product (3.22) has pure point spectrum iff $\alpha(s)$ is not cohomologous to a constant.

(3) A constant function $\alpha(s) = \alpha$ is a coboundary iff α and γ are rationally dependent.

(4) By (1-3) for a function $\alpha(s)$ cohomologous to a constant α the skew product is either not ergodic or ergodic with pure point spectrum, depending on whether α and γ are rationally dependent or not.

The following result is a consequence of Theorem 1 from [BM] combined with Lemma 3.12.

Lemma 3.13. Consider an Anzai skew product (3.22) with γ Diophantine and $\alpha : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ of zero degree and C^∞ . Then there exists a number $\alpha \in [0, 1)$ such that the map F is smoothly (i.e. C^∞) conjugate to

$$(3.24) \quad F_0(\mathbb{S}^1 \times \mathbb{S}^1) \rightarrow (\mathbb{S}^1 \times \mathbb{S}^1), \quad F_0(s, \varphi) = (s + \gamma, \varphi + \alpha).$$

We prepare for our main result related to the full system, Theorem 3.15, with the proof of the following Lemma.

Lemma 3.14. (1) The map $\alpha(s)$ is of zero degree, thus can be regarded as $\alpha : \mathbb{S}^1 \rightarrow \mathbb{R}$.

(2) Both $\alpha(s)$ and $\tau(s)$ are C^∞ .

Proof. (1) Note that any continuous map $\mathbb{S}^1 \rightarrow \mathbb{S}^1$ that has nonzero degree must be onto. On the other hand, as we are in the alternating case we have $\cos \hat{\beta}(s) > R$ for any $s \in \mathbb{S}^1$ (see Figure 1). This means

$|\alpha(s)| \leq 2\arccos R$ which in particular implies $\alpha(s)$ cannot be onto, thus it is of zero degree.

(2) The maps $\alpha(s)$ and $\tau(s)$ are defined in Formulas (3.21) and (3.23) respectively (see also (3.15)). To establish the desired smoothness properties we need to show that four functions, namely (i) $\beta(s)$ (ii) $\vec{v}(s)$ (iii) $\hat{\beta}(\beta)$ (iv) $d(\beta)$ are all smooth. Moreover we need to show that (v) $|\vec{v}(s)|$ is bounded away from zero.

First note x, y, z all depend on s smoothly, while w a priori only continuously. To see this note that $x(s), y(s), z(s)$ are the Cartesian coordinates of a circle (of the velocity circle, cf. Section 3.2) thus depend smoothly on the arclength parameter. By (3.10) $w(s)$ is the square root of a smooth non-negative function, thus it is continuous, and smooth if bounded away from zero, a fact we establish below.

By Remark 3.1 (2) and compactness $|\vec{v}(s)|$ is bounded away from zero, thus we have statement (v). This by alternating (cf. (3.17)) gives positive lower bound and, by the above reasoning, smoothness for $w(s)$ which means (cf. (3.8)) we have shown (ii). Moreover (i) follows as we have demonstrated that $\tan \beta(s)$ is smooth.

We need to prove (iii) and (iv). As we are in the alternating case, see the first part of Remark 3.6, we have Formula (3.17) and thus by compactness

$$R^2 - \sin^2 \beta(s) > K_1 > 0; \quad \text{and} \quad \frac{\sin^2 \beta}{R^2} > K_2 > 1.$$

This means that the quantity under the square root in the formula for $d(\beta)$ (cf. (3.23)) is bounded away from zero. Similarly, in the formula for $\hat{\beta}(\beta)$ (see (3.21)), the argument of arc sin is bounded away from 1 and -1 . Thus we have (iii) and (iv). \square

Theorem 3.15. *Let us fix γ Diophantine.*

- (1) *The map (3.22) is smoothly conjugate to (3.24) with $\alpha = \int_{\mathbb{S}^1} \alpha(s) ds$.*
- (2) *There is a positive measure set of integrals of motion for which (3.22) is uniquely ergodic.*
- (3) *The system in “physical” time, i.e. the suspension flow over (3.22) with ceiling function (3.23) is never ergodic.*

Proof. (1) Lemmas 3.14 and 3.13 imply that (3.22) is smoothly conjugate to (3.24).

(2) For α and γ rationally independent the system (3.24) is well-known to be uniquely ergodic, see [KH]. What we are going to demonstrate is that this happens for a positive measure set of parameters. To see this, first note that the velocity circle is determined by the two parameters N and E (see Equation (3.9) and the discussion at the beginning of Section 3.2). Furthermore, it is possible to choose these two parameters from a positive measure set in \mathbb{R}^2 to ensure that

$x(s) = v_t(s) > 0$ for all $s \in \mathbb{S}^1$. We keep the physical parameters (and thus γ) together with two integrals of motion, with N and E from the above set fixed and vary the third one, F from (3.10). What is shown below is that in this setting α depends on F in a continuous and strictly monotonic manner. This implies that α and γ can be rationally dependent only for countably many values of F , and thus by our choice of N and E for a positive measure set of integrals.

In the rest of the proof, all functions of s (e.g. $z(s), \beta(s), \alpha(s)$) refer to the original values of the integrals, while the primed ones (e.g. $z'(s), \beta'(s), \alpha'(s)$) indicate maps for F slightly increased. Since Lebesgue measure on the circle is invariant with respect to the transformation $s \rightarrow s_2$, we have

$$\alpha = \int_{\mathbb{S}^1} (\hat{\beta}(s) + \hat{\beta}(s_2)) ds = 2 \int_{\mathbb{S}^1} \hat{\beta}(s) ds.$$

Note that $x(s) = x'(s)$ and $y(s), z(s)$ similarly, however as F is increased we get $w'(s) > w(s)$ for all $s \in \mathbb{S}^1$. This implies $\beta'(s) < \beta(s)$ (recall our convention on $x(s) > 0$) which yields, by a simple geometrical calculation, $\hat{\beta}'(s) < \hat{\beta}(s)$ for all $s \in \mathbb{S}^1$. Thus α depends on F strictly monotonically.

(3) A suspension flow above an ergodic map is ergodic iff the ceiling function is not cohomologous to a constant, see [KH]. As $\tau(s)$ is C^∞ by Lemma 3.14 and γ is chosen to be Diophantine, we may apply Lemma 3.13 (note also Definition 3.10) to see that $\tau(s)$ is smoothly cohomologous to its average, a constant, thus the suspension flow is not ergodic. \square

Remark 3.16. Note that an Anzai skew product (3.20) is never mixing as it has a factor which is a rotation by γ (with the semi-conjugacy being projection to the first \mathbb{S}^1 -coordinate). On the other hand, in case the spectrum is not pure point only “horizontal” sets $B' = B \times \mathbb{S}^1$ do not mix weakly, see Lemma 3.12. To be more precise, if the spectrum is not pure point, restricting the unitary transformation of $\mathcal{L}^2(\mathbb{S}^1 \times \mathbb{S}^1)$ to the orthocomplement of the subspace generated by characteristic functions of horizontal sets, the spectrum is continuous (see [A]). Even though this phenomenon, if occurs at all in our system, can only occur for a zero measure set of parameters (in the non-Diophantine case), it seems to be an interesting point for further investigation.

4. TWO PARTICLES, ONE ROTATING SCATTERER

In this final section we will show that the motion of two particles with one rotating scatterer, although more complicated, resembles the motion of one particle with two rotating scatterers. Let the two point particle velocity vectors be \vec{v} and \vec{u} . Note that as we assume that only the inner scatterer rotates, the outer interactions are billiard collisions

and thus for both particles the argument of Lemma 2.1, statement (1) applies. This means that both of the particles either alternate (all odd bounces are in and all even ones are out) or collide always with the outer billiard scatterer. We may assume that both particles alternate. Otherwise motion is a direct product of two integrable systems: two circle billiards if both particles remain outside and a product of a circle billiard with the system discussed in Section 2 if one of them alternates.

Recall Remark 1.1: in the present section splitting of \vec{v} (\vec{u}) into tangential and normal components (v_t, v_n) ((u_t, u_n)) is always understood at the *inner* scatterer. Again by the geometrical picture discussed in the proof of Lemma 2.1 both u_n and v_n are integrals of motion. For brevity we drop subscripts and denote v_t (u_t) by v (u).

We need to introduce one more velocity coordinate ω for the angular velocity of the inner scatterer. Furthermore, as in Section 2, $\eta = \frac{\Theta}{mR^2}$ where Θ is the moment of inertia of the scatterer and m is the mass of one of the point particles (we assume them to have the same mass).

An interaction between the particle with velocity \vec{v} and the inner scatterer will be referred to as a v-collision, and the definition of u-collision is analogous. In course of a v-collision the velocities transform as

$$(4.1) \quad \begin{aligned} v' &= v - \frac{2\eta}{1+\eta}(v - R\omega); \\ R\omega' &= R\omega + \frac{2}{1+\eta}(v - R\omega). \end{aligned}$$

We have

$$(4.2) \quad v + \eta R\omega = v' + \eta R\omega'; \quad v^2 + \eta R^2\omega^2 = v'^2 + \eta R^2\omega'^2$$

while u does not change. For u-collisions

$$(4.3) \quad \begin{aligned} u' &= u - \frac{2\eta}{1+\eta}(u - R\omega); \\ R\omega' &= R\omega + \frac{2}{1+\eta}(u - R\omega), \end{aligned}$$

and we have conserved quantities:

$$(4.4) \quad u + \eta R\omega = u' + \eta R\omega'; \quad u^2 + \eta R^2\omega^2 = u'^2 + \eta R^2\omega'^2$$

in addition to the unchanged v .

Let us introduce $q = \sqrt{\eta}R\omega$. Formulas (4.2) and (4.4) imply the presence of two integrals:

$$v^2 + u^2 + q^2 = E \quad \text{and} \quad v + u + \sqrt{\eta}q = N.$$

Thus velocity motion is restricted to the intersection of a sphere and a plane in \mathbb{R}^3 , a circle which will be referred to as the velocity circle.

The v- and u-collisions, i.e. the transformations (4.1) and (4.3) are reflections across the planes

$$\sqrt{\eta}v - q = 0 \quad \text{and} \quad \sqrt{\eta}u - q = 0,$$

respectively. When restricted to the velocity circle, v-collisions reduce to reflections across the line

$$(4.5) \quad (N + t, -(1 + \eta)t, \sqrt{\eta}t)$$

while u-collisions to reflections across

$$(- (1 + \eta)t, N + t, \sqrt{\eta}t).$$

The angle of the two lines, which we denote by $\frac{\gamma}{2}$, is an important parameter depending on which we distinguish rational/irrational cases. We may calculate

$$\cos \frac{\gamma}{2} = \frac{1}{1 + \eta},$$

thus, just like in Lemma 3.2, for a full measure set of η parameters – actually, apart from countably many points – the base is irrational, and this classification does not depend on the integrals of motion.

We parameterize the velocity circle with the normalized arclength parameter $s \in [0, 1]$ as in Convention 3.3. For definiteness $s = 0$ corresponds to one of the intersection points of the circle with the line (4.5). In this coordinate the v- and u-collisions, i.e. the transformations (4.1) and (4.3) are

$$(4.6) \quad s \rightarrow s_1 = -s \pmod{1} \text{ and}$$

$$(4.7) \quad s \rightarrow s_2 = -s - \gamma \pmod{1},$$

respectively.

However, at this point the exposition ceases to be parallel to that of Section 3.2. The two transformations above may follow each other in a more complicated manner and it is not enough to take two copies of the circle to describe the velocity evolution.

Consider the Poincaré section for the cross section of v-collisions. When the system is in state s , there is a unique tangential velocity $v(s)$ – and thus, as the normal velocity component v_n is an integral of motion, a unique velocity vector $\vec{v}(s)$ – for the first particle. Thus the return time for the first particle can be calculated (cf. Fig 2) as

$$(4.8) \quad \tau_1(s) = \tau(v(s)), \quad \text{where} \quad \tau(v) = 2 \frac{d(\vec{v})}{|\vec{v}|}$$

and $d(\vec{v})$ is given in Formula (3.23) (note this time v_n does not depend on s as it is an integral of motion). The evolution of the system strongly depends on the number of inner bounces the second particle has between two consecutive bounces of the first one. As within this time interval the second particle follows the pattern described in Section 2, we have the following picture. Let us assume the system is

in state $s \in \mathbb{S}^1$ just after a v-collision and thus the angular velocity of the scatterer is $\omega(s)$. If the number of u-collisions before the next v-collision is even, the scatterer, just before the next v-collision, will rotate with the exact same angular velocity $\omega(s)$. This implies that just after the second v-collision the system arrives at the state $s_1 = -s$ (cf. (4.6)). On the other hand, if the number of u-collisions in this time interval is odd, the scatterer rotates with angular velocity $\omega(s_2)$ just before the second v-collision, and thus the system lands just after the second v-collision at the state $-s_2 = s + \gamma$ (cf. (4.7)).

To calculate the number of u-collisions within the time interval $[0, \tau_1]$ we denote by

$$(4.9) \quad \tau_2(s) = \tau(u(s)) \quad \text{and} \quad \tau_3(s) = \tau(u(s_2))$$

the length of time intervals between consecutive u-collisions in the two possible phases of the second particle (cf. also (4.8)). We introduce

$$(4.10) \quad \hat{\lambda} = \frac{\tau_2}{\tau_2 + \tau_3} \in [0, 1) \quad \text{and} \quad \hat{t} = \left\{ \frac{\tau_1}{\tau_2 + \tau_3} \right\} \in [0, 1)$$

(where $\{x\}$ is the fractional part of $x \in \mathbb{R}$). If the two particles started from the inner scatterer at the same time, the parity of u-collisions until the next v-collision would depend only on the relation of these two quantities: it would be even/odd for \hat{t} less than (greater than) $\hat{\lambda}$. To track the velocity evolution in the general case we need to introduce one more coordinate.

Definition 4.1. *At the moment of a v-collision (i.e. on our Poincaré section) we define $t \in [0, 1)$ as the following quantity: the last u-collision preceded the v-collision we are considering at time $t\tau_2(s)$.*

The coordinate t defined this way describes where the second particle is at the moment the first particle collides inside.

Convention 4.2. *We think of the quantities t, \hat{t} as elements of \mathbb{S}^1 (identified with $[0, 1)$) thus we may add them (which is always understood (mod 1)).*

This way the velocity evolution can be described as the following two dimensional system

$$(4.11) \quad \begin{aligned} T : \mathbb{S}^1 \times \mathbb{S}^1 &\rightarrow \mathbb{S}^1 \times \mathbb{S}^1, \quad T(s, t) = (s', t') \quad \text{where} \\ (s', t') &= \left(-s, t + \frac{\hat{t}}{\hat{\lambda}} \right) \quad \text{if} \quad t\hat{\lambda} + \hat{t} < \hat{\lambda} \quad (\text{even case}) \\ (s', t') &= \left(s + \gamma, \frac{\tau_2}{\tau_3}t + \frac{\hat{t}}{1 - \hat{\lambda}} \right) \quad \text{if} \quad t\hat{\lambda} + \hat{t} > \hat{\lambda} \quad (\text{odd case}). \end{aligned}$$

Note that T is well defined since \hat{t} and $\hat{\lambda}$ both depend on s , see Equation (4.10).

Remark 4.3. *The reason for the rescaling term $\frac{\tau_2(s)}{\tau_3(s)}$ in the evolution of the t -coordinate for the odd case is that, as there has been an odd number of u -collisions, the length of the period for the second particle is τ_3 in contrast to the previous τ_2 .*

We have $t' = \frac{t\tau_2(s) + \hat{t}(\tau_2(s) + \tau_3(s))}{\tau_2(s')}$ for both the odd and the even cases. Note $u(s) = u(-s)$ (as the reflection (4.6) corresponds to the interaction of the first particle with the scatterer, thus the velocity of the second particle does not change) and thus $\tau_2(s) = \tau_2(-s)$ (cf. (4.8) and (4.9)). This implies $\tau_2(-s) = \tau_2(s)$ (even case) and $\tau_2(s + \gamma) = \tau_3(s)$ (odd case).

Note that this way we need to study a two dimensional system for velocity evolution, which is much more complicated than the one dimensional base system of Section 3.2. However, as the theorem below shows, for a suitable choice of integrals and, correspondingly, initial conditions (both of positive measure), there is an arbitrary long time interval for which the t coordinate can be disregarded and for which the velocity motion follows a rotation orbit.

Theorem 4.4. *For any $\varepsilon > 0$ small enough there exists a positive integer N_ε , an open set W in the space of parameters (i.e. integrals of motion) and a set of initial conditions E of Lebesgue measure bigger than ε with the following property. Given a system with parameters from W , for any $(s_0, t_0) \in E$ the iterate $(s_k, t_k) = T^k(s_0, t_0)$ is in the odd part of the phase space, and thus, $s_k = s_0 + k\gamma$, as long as $0 \leq k < N_\varepsilon$. Moreover, $N_\varepsilon \rightarrow \infty$ as $\varepsilon \rightarrow 0$.*

Proof. Let us define $E \subset \mathbb{S}^1 \times \mathbb{S}^1$ with $|t_0 - \frac{1}{2}| < \varepsilon$ and $s_0 \in \mathbb{S}^1$ arbitrary. This means that the second particle is close to the outer (billiard) scatterer when the first particle collides with the inner (rotating) one.

To define the open set W let us fix N and E arbitrary (in such a way that the velocity circle is not empty). Note that the normal velocity components are integrals of motion as well. We fix $\varepsilon' > 0$ small enough and $K > 0$ big enough, their values will be chosen later depending on ε . For the points of W we assume $|v_n - K'| < \varepsilon'$ and $|u_n - K'| < \varepsilon'$ for some $K' > K$. This ensures that the normal velocity components of the particles are almost equal and, moreover, much greater than the tangential components. As a consequence, for all s , both $d(v(s))$ and $d(u(s))$ are very close to $1 - R$, while both $|\vec{v}(s)|$ and $|\vec{u}(s)|$ are very close to K' . In particular, with ε' sufficiently small and K sufficiently big we may ensure that

$$(4.12) \quad \begin{aligned} \hat{t}(s) &\in \left(\frac{1}{2} - \varepsilon, \frac{1}{2} + \varepsilon \right), & \hat{\lambda}(s) &\in \left(\frac{1}{2} - \varepsilon, \frac{1}{2} + \varepsilon \right), \\ \frac{\tau_2(s)}{\tau_3(s)} &\in (1 - \varepsilon, 1 + \varepsilon), & \text{and} & \left| \frac{\hat{t}(s)}{1 - \hat{\lambda}(s)} \right| < \varepsilon \end{aligned}$$

for all $s \in \mathbb{S}^1$, cf. (4.10). Note that in the last inequality the absolute value is understood as distance from 0 on \mathbb{S}^1 .

Finally, we fix

$$(4.13) \quad N_\varepsilon = \left[\min \left(\frac{\log(2 - 2\varepsilon)}{\log(1 + 3\varepsilon)}, \frac{\log(8\varepsilon)}{\log(1 - 3\varepsilon)} \right) \right] - 2,$$

where $[x]$ denotes the integer part of $x \in \mathbb{R}$.

Lemma 4.5. *Assume $0 \leq k < N_\varepsilon$. We have*

$$(1) \quad \frac{1}{2}(1 - 3\varepsilon)^{k+1} < t_k < \frac{1}{2}(1 + 3\varepsilon)^{k+1}, \quad (2) \quad \hat{\lambda} < t_k \hat{\lambda} + \hat{t} < 1.$$

Note that (2) means in particular that (s_k, t_k) is in the odd part of the phase space, cf. (4.11).

Proof. We will prove (1) and (2) by induction. As $|t_0 - \frac{1}{2}| < \varepsilon$, both statements hold for $k = 0$ by (4.12). Assume that we have already established (1) and (2) for $k - 1$. Note that this, in particular, means that (s_{k-1}, t_{k-1}) is in the odd part and thus by (4.11)

$$t_k = \frac{\tau_2}{\tau_3} t_{k-1} + \frac{\hat{t}}{1 - \hat{\lambda}}.$$

This implies, by the inductive assumption and (4.12)

$$\begin{aligned} t_k &< \frac{1}{2}(1 + 3\varepsilon)^k(1 + \varepsilon) + \varepsilon &< \frac{1}{2}(1 + 3\varepsilon)^{k+1}, \\ t_k &> \frac{1}{2}(1 - 3\varepsilon)^k(1 - \varepsilon) - \varepsilon &> \frac{1}{2}(1 - 3\varepsilon)^{k+1} \end{aligned}$$

which is (1) for t_k . On the other hand, (1) implies (2) as long as $k < N_\varepsilon$. To see this note that by (4.12) it is enough to prove that

$$\begin{aligned} \frac{1}{2}(1 - 3\varepsilon)^{k+1} \frac{1}{2}(1 - 2\varepsilon) + \frac{1}{2} - \varepsilon &> \frac{1}{2} + \varepsilon \quad \text{and} \\ \frac{1}{2}(1 + 3\varepsilon)^{k+1} \frac{1}{2}(1 + 2\varepsilon) + \frac{1}{2} + \varepsilon &< 1. \end{aligned}$$

These conditions are both satisfied as long as $k < N_\varepsilon$, cf. (4.13). \square

By Formula (4.11), Theorem 4.4 is a straightforward consequence of Lemma 4.5. \square

Remark 4.6. *Considering motion in the fibers, i.e. in addition to the velocity coordinate s , the consecutive points of impact of v -collisions (points of impact for the first particle) parametrized by the arclength parameter φ along the inner scatterer, we can make the following observation. Under the conditions of Theorem 4.4, the coordinate t can be disregarded, and the two dimensional system of the (s, φ) variables follows (up to time N_ε) the Anzai skew product dynamics analyzed in Section 3.3.*

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